Self-adjoint codimension 2 boundary conditions for Dirac operators

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Section: Differential Geometry and Global Analysis Organized by: Vicente Cortés and Oliver Goertsches Annual Meeting of the DMV in Chemnitz (virtual) Sept 15, 2020

Slides available on

http://www.mathematik.uni-regensburg.de/
 ammann/talks/2020DMV-handout.pdf



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Classical boundary conditions for the Laplace operator

Let  $\Omega \subset \mathbb{R}^n$  be a bounded domain with smooth boundary. To model problems in physics or technical applications we often want to solve the Poisson problem

$$\Delta u = f$$

for a given function  $f \in C^{\infty}(\Omega)$  or the eigenvalue problem

$$\Delta u = \lambda u$$

for  $\lambda \in \mathbb{R}$ .

**Many** solutions, and many solutions behave unphysically at the boundary.

~ We need boundary conditions!



# The minimal and the maximal operator We define

 $\Delta_{\boldsymbol{c}}: \boldsymbol{C}^{\infty}_{\boldsymbol{c}}(\Omega) 
ightarrow \boldsymbol{C}^{\infty}_{\boldsymbol{c}}(\Omega)$ 

as a densely defined operator in  $L^2(\Omega)$ .

The minimal operator  $\Delta_{\min}$  is the closure of  $\Delta_c$  with respect to the graph norm.

$$L^{2}(\Omega) \supset H^{2}_{0}(\Omega) \xrightarrow{\Delta_{\min}} L^{2}(\Omega)$$

And the maximal operator is defined as its adjoint

$$\underbrace{\left\{ \begin{matrix} u \in L^2(\Omega) \mid \Delta u \in L^2(\Omega) \\ \\ \vdots = \mathsf{dom}(\Delta_{\max}) \supsetneq H^2(\Omega) \end{matrix} \right\}}_{:= \mathsf{dom}(\Delta_{\max}) \supsetneq H^2(\Omega)} \xrightarrow{\Delta_{\max}} L^2(\Omega).$$

The map

$$\begin{aligned} \tau = (\tau_{\mathcal{D}}, \tau_{\mathcal{N}}) : \mathcal{C}^{\infty}(\overline{\Omega}) &\to \mathcal{C}^{\infty}(\partial\Omega) \times \mathcal{C}^{\infty}(\partial\Omega), \\ u &\mapsto (u|_{\partial\Omega}, (\partial_{\nu}u)|_{\partial\Omega}) \end{aligned}$$

extends to a trace map  $\tau$  : dom $(\Delta_{\max}) \rightarrow \check{H}(\partial \Omega)$ .



#### Symplectic structure on boundary data

Trace map  $\tau : dom(\Delta_{max}) \to \check{H}(\partial\Omega)$ . Here  $\check{H}(\partial\Omega)$  is a Sobolev type Hilbert space,

$$H^{3/2}(\partial\Omega) imes H^{1/2}(\partial\Omega) \subset \check{H}(\partial\Omega) \subset H^{-1/2}(\partial\Omega) imes H^{-3/2}(\partial\Omega)$$

on which a non-degenerate perfect symplectic pairing

$$\begin{split} \check{H}(\partial\Omega) \times \check{H}(\partial\Omega) & \xrightarrow{B} \mathbb{R} \\ \left( (v_1, w_1), (v_2, w_2) \right) & \mapsto \int_{\partial\Omega} (v_1 w_2 - v_2 w_1) \end{split}$$

is well-defined. Green identity

$$\int_{\Omega} (\Delta u_1) u_2 - \int_{\Omega} u_1 \Delta u_2 = \pm B(\tau(u_1), \tau(u_2)).$$



## Self-adjoint extensions

If  $A \subset \check{H}(\partial \Omega)$  is a closed subspace, then

$$\underbrace{\tau^{-1}(A)}_{\subset \operatorname{dom}(\Delta_{\max})} \xrightarrow{\Delta_A := \Delta_{\max}|_{\tau^{-1}(A)}} L^2(\Omega)$$

is a closed extension.

 $\Delta_A$  is self-adjoint iff A is a Lagrangian subspace of  $(\check{H}(\partial\Omega), B)$ .

#### Examples:

Dirichlet boundary conditions:  $A := \{(0, w) \in \check{H}(\partial \Omega)\}$ Neumann boundary conditions:  $A := \{(v, 0) \in \check{H}(\partial \Omega)\}$ 

Under some regularity assumptions:

- solution of the Poisson problem
- discrete, real spectrum
- eigenspaces finite-dimensional, spanned by smooth functions



## Codimension 1 bdy cond. for Dirac operators

#### Now replace:

$\overline{\Omega} \subset \mathbb{R}^n$	$\mathcal{C}^{\infty}(\overline{\Omega})$	Laplacian $\Delta$
cpct. Riem. spin manfd.	twisted spinors	twisted Dirac oper.
with boundary <i>M</i>	$\mathcal{C}^{\infty}(M, V \otimes \Sigma M)$	Ø

$$\begin{split} \tau =: \mathbf{C}^{\infty}(\mathbf{M}, \mathbf{V} \otimes \Sigma \mathbf{M}) &\to \mathbf{C}^{\infty}\big(\partial \mathbf{M}, (\mathbf{V} \otimes \Sigma \mathbf{M})|_{\partial \mathbf{M}}\big), \\ \mathbf{u} &\mapsto \mathbf{u}|_{\partial \Omega} \end{split}$$

extends to a trace map  $\tau : \operatorname{dom}(\mathcal{D}_{\max}) \to \check{H}(\partial M)$ . Here  $\check{H}(\partial \Omega)$  is a Sobolev type Hilbert space,

$$H^{1/2}\big(\partial M, (V\otimes \Sigma M)|_{\partial M}\big)\subset \check{H}(\partial M)\subset H^{-1/2}\big(\partial M, (V\otimes \Sigma M)|_{\partial M}\big)$$

with a skew-hermitian sesquilinear map  $B: \check{H}(\partial M) \times \check{H}(\partial M) \rightarrow \mathbb{C}.$ 

$$\int_{\boldsymbol{M}} \langle \boldsymbol{D} \varphi, \psi \rangle - \int_{\boldsymbol{M}} \langle \varphi, \boldsymbol{D} \psi \rangle = \boldsymbol{B} \big( \tau(\varphi), \tau(\psi) \big).$$



## Self-adjoint extensions for D

If  $A \subset \check{H}(\partial M)$  is a closed subspace, then

$$\underbrace{\tau^{-1}(A)}_{\text{Ldom}(\emptyset_{\max})} \xrightarrow{\emptyset_A := \emptyset_{\max}|_{\tau^{-1}(A)}} L^2(M, V \otimes \Sigma M)$$

is a closed extension.

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 $\mathcal{D}_A$  is self-adjoint iff A is a Lagrangian subspace of  $(\check{H}(\partial M), B)$ .

Under some regularity assumptions:

- discrete, real spectrum
- eigenspaces finite-dimensional, spanned by smooth sections
- $D_A$  is a Fredholm operator

→ Atiyah-Patodi-Singer index theorem. See e.g. Bär-Ballmann arxiv 1101.1196.



## The setting of our work

- Let (M, g) be a complete oriented Riemannian manifold, N a compact oriented submanifold of codimension 2.
- We assume that  $M \setminus N$  is spin. Thus there is a complex spinor bundle  $\Sigma \to M \setminus N$ .
- Let  $L \to M \setminus N$  be a flat hermitian line bundle. We get

$$\pi_1(M \setminus N) \to S^1 \subset \mathbb{C}$$

•  $W := \Sigma \otimes L$  generalized spinor bundle over  $M \setminus N$ More general frameworks are possible which will not be discussed in this talk.



### Monodromy

Monodromy  $\alpha = (\alpha_1, \ldots, \alpha_j)$ .

$$N = \prod_{j=1}^{\ell} N_j$$

decomposition into connected components

Parallel transport in *W* around  $N_j$  is  $e^{2\pi i \alpha_j}$ .  $[\alpha_j] \in \mathbb{R}/\mathbb{Z}$  only depends on *j*.



#### Main examples

M spin. Monodromy comes from L. Main subcase: L flat. Monodromy π<sub>1</sub>(M \ N) → S<sup>1</sup>. Main subsubcase: N is a link in S<sup>3</sup>.

 $(S^1)^\ell \ni (\exp 2\pi i \alpha_j)_{j \in \{1,...,\ell\}} \mapsto L_\alpha$ 

*M* is not spin, (more precisely: spin structure does not extend), *N* connected. Then monodromy only comes from Σ, α = 1/2 mod Z. Main subcase: L = C Example: M = CP<sup>2r</sup>, N = CP<sup>2r-1</sup>. Fix p ∈ M \ N, solve ØΨ = ψ<sub>0</sub>δ<sub>p</sub> on M \ N with bdy cond. Expectation: If PMT would fail, we would get a map

 $S(\Sigma_{\rho}) \times \{ bdy \ cond \} \rightarrow \{ non-zero \ spinors \ on \ N \}.$ 

Interesting applications?



#### Results by Ammann and Große

We obtain

• a bundle  $S \rightarrow N$ ,

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- a Hilbert space of sections  $\check{H}(N, S)$ ,
- a skew-hermitian sesquilinear non-degenerate perfect pairing B : H̃(N, S) × H̃(N, S) → C
- ▶ a trace map  $\tau$  : dom( $D_{max}$ ) →  $\check{H}(N, S)$

If  $A \subset \check{H}(N, S)$  is a closed subspace, then

$$\underbrace{\tau^{-1}(A)}_{\text{Ldom}(\emptyset_{\max})} \xrightarrow{\emptyset_A := \emptyset_{\max}|_{\tau^{-1}(A)}} L^2(M, V \otimes \Sigma M)$$

is a closed extension.  $\operatorname{dom}(\mathcal{D}_{\emptyset}) = \operatorname{dom}(\mathcal{D}_{\min}) \text{ and } \operatorname{dom}(\mathcal{D}_{\check{H}(N,S)}) = \operatorname{dom}(\mathcal{D}_{\max}).$  $\mathcal{D}_A$  is self-adjoint iff A is a Lagrangian subspace of  $(\check{H}(\partial M), B).$ 

#### Some remarks

▶ In the case  $\alpha \in \mathbb{Z}$  the submanifold *N* is invisible, i.e.

$$ot\!\!/ p_{\min} = 
ot\!\!/ p_{\max} = 
ot\!\!/ p^M$$

►  $\tau$  is not the extension of a restriction map. In contrast we have: Suppose that  $\varphi \in \text{dom}(D_{\text{max}})$  is bounded on a neighbourhood of *N*. Then  $\varphi \in \text{dom}(D_{\text{min}})$ .

• The bundle  $S \rightarrow N$  has a Clifford multiplication

$$TM|_N \otimes S \to S$$

Normal volume element:  $\omega_{nor} := e_1 \cdot e_2$  if  $(e_1, e_2)$  is a positively oriented orthonormal basis of the normal bundle. This gives a splitting

$$\mathcal{S} = \mathcal{S}^+ \oplus \mathcal{S}^-$$

into the  $\pm i$ -eigenspace bundles for the Clifford action of  $\omega_{nor}$ .



#### Example: Portmann & Sok & Solovej

 $\check{H}(N, S^+)$  and  $\check{H}(N, S^-)$  are Lagrangian subspaces of  $\check{H}(N, S)$ .

PSS (2015–2018) considered the special case that  $M = S^3$  is the round sphere and N is a link.

Electrons coupled magnetic fields.

Existence of harmonic spin<sup>*c*</sup>-spinors yield statements of the type

If our world is stable, then the fine structure constant  $\hbar c/e^2$  has to satisfy some bounds.

Measurements:  $\hbar c/e^2 = 137.03599968...$ 

A spectral flow argument yields harmonic spinors. Question: How is this spectral flow related to classical link invariants?



#### 2-dimensional model space

Assume  $M = \mathbb{C} \ni z$ ,  $N = \{0\}$ ,  $\Sigma = \underline{\mathbb{C}^2} = \Sigma_+ \oplus \Sigma_-$ *L* flat bundle over  $[\mathbb{C} : \{0\}]$ , monodromy  $\alpha$ Then  $\omega_{nor}$  is the standard volume element.

 $\frac{z^{-\alpha}}{|z|^{-\alpha}}$  represents a nowhere vanishing smooth section of *L*. Ansatz:

$$\Phi^+_{\beta,\gamma} := \begin{pmatrix} z^\beta \overline{z}^\gamma \\ 0 \end{pmatrix}, \qquad \Phi^-_{\beta,\gamma} := \begin{pmatrix} 0 \\ z^\beta \overline{z}^\gamma \end{pmatrix}.$$

where  $\beta$  and  $\gamma$  over real numbers with  $\beta - \gamma + \alpha \in \mathbb{Z}$ .  $\Phi_{\beta,\gamma}^{\pm} \in \mathcal{L}_{\text{loc}}^{2}$  iff  $\beta + \gamma > -1$ 

$$\not\!\!\!D \Phi^+_{\beta,\gamma} = -\sqrt{2}\beta \Phi^-_{\beta-1,\gamma}, \qquad \not\!\!\!D \Phi^-_{\beta,\gamma} = \sqrt{2}\gamma \Phi^+_{\beta,\gamma-1},$$



#### Lemma.

The condition that  $\Phi_{\beta,\gamma}^{\pm} \in \text{dom}(\mathcal{D}_{\max})$  is characterized as follows ("locally around 0").

- (1) Suppose  $\beta \neq 0$  and  $\gamma \neq 0$ . Then  $\Phi_{\beta,\gamma}^{\pm} \in \text{dom}(\not\!\!D_{\text{max}})$  if and only  $\beta + \gamma > 0$ .
- (2) Suppose  $\beta = 0$  and  $\gamma \neq 0$ . Then  $\Phi_{0,\gamma}^+ \in \text{dom}(\not\!\!D_{\text{max}})$  if and only if  $\gamma > -1$ , and  $\Phi_{0,\gamma}^- \in \text{dom}(\not\!\!D_{\text{max}})$  if and only if  $\gamma > 0$ .
- (3) Suppose  $\beta \neq 0$  and  $\gamma = 0$ . Then  $\Phi_{\beta,0}^+ \in \text{dom}(\not\!\!D_{\text{max}})$  if and only if  $\beta > 0$ , and  $\Phi_{\beta,0}^- \in \text{dom}(\not\!\!D_{\text{max}})$  if and only if  $\beta > -1$ .

(4) Suppose 
$$\beta = \gamma = 0$$
.  $\Phi_{0,0}^{\pm} \in \operatorname{dom}(\mathcal{D}_{\max}) = \operatorname{dom}(\mathcal{D}_{\min})$ .

 $lpha \in (0,1)$ : Then elements in dom $({
otin max})$  are of the form

$$\begin{pmatrix} \overline{z}^{\alpha-1}\varphi_+\\ z^{-\alpha}\varphi_- \end{pmatrix} + \operatorname{dom}(\not\!\!D_{\min}).$$



#### Disclaimer

The project is still work in progress. Until now we have only written up in detail the case of a totally geodesic submanifold N and some similar assumptions, but we do not expect any modifications for the general case. Preprints are not yet available.

#### Summary

- We have a complete description of the self-adjoint extensions of twisted Dirac operators on manifolds M<sup>m</sup> \ N<sup>m-2</sup>.
- ▶ The space *N* is invisible, if the monodromy is trivial.
- Spinors in the domain of the maximal operator which are not in the minimal domain are not bounded (near the boundary). The trace map has to be modified considerably.
- Boundary data are given by elements in a Hilbert space H̃ of sections of a bundle S → N. This space depends strongly on the monodromy.

#### Thanks...

... to Boris Botvinnik and Nikolai Saveliev for discussions about link invariants associated to these results.

